THE FINITE ELEMENT METHOD 2013 Dept. of Solid Mechanics

Short solutions: 2013-06-05

Problem 1:

a) Multiply

$$-\gamma\dot{\rho} = \rho + \gamma div(\boldsymbol{\nabla}\rho) - \lambda$$

by a scalar weight function, v, and integrate over the domain, V. Use Green-Gauss theorem on the divergence term. Split the boundary into $S = S_c + S_h$. The result is:

$$-\int_{V} v\gamma \dot{\rho} dV = \int_{V} v\rho dV - \int_{V} \gamma \left(\nabla v\right)^{T} \nabla \rho dV + \int_{S_{c}} v(1+\rho) dS - \int_{V} v\lambda dV$$

b) Use of the interpolation, $\rho = N\rho$ yields $\dot{\rho} = N\dot{\rho}$ and $\nabla \rho = B\rho$. The weight function, v, is chosen as v = Nc which yields $\nabla c = Bc$. Insertion into the weak form results in

$$\int_{V} \gamma \boldsymbol{B}^{T} \boldsymbol{B} dV \boldsymbol{\rho} - \int_{V} \boldsymbol{N}^{T} \boldsymbol{N} dV \boldsymbol{\rho} - \int_{V} \boldsymbol{N}^{T} \gamma \boldsymbol{N} dV \dot{\boldsymbol{\rho}} - \int_{S_{c}} \boldsymbol{N}^{T} \boldsymbol{N} \boldsymbol{\rho} dS = \int_{S_{c}} \boldsymbol{N}^{T} dS - \int_{V} \boldsymbol{N}^{T} \lambda dV$$

Problem 2 :

The contribution to node 1 is given as

$$\{\boldsymbol{f}_b\}_1 = \int_{\mathcal{L}} N_1 \boldsymbol{t} d\mathcal{L} = \int_{\mathcal{L}_{right}} N \boldsymbol{t} d\mathcal{L} + \int_{\mathcal{L}_{left}} N_1 \boldsymbol{t} d\mathcal{L} = [symmetry] = 2 \int_{\mathcal{L}_{right}} N_1 \boldsymbol{t} d\mathcal{L}$$

In the element located to the right, the shape function associated with node 1 can be expressed as $N_1^e = (a - \eta)(2a - \eta)(3a - \eta)/(6a^3)$

$$2\int_{\mathcal{L}_{right}} N_1 t d\mathcal{L} = 2\int_0^{3a} N_1 t d\eta = 2\int_{-1}^1 N_1 t \frac{d\eta}{d\xi} d\xi =$$
$$2\left(\left(N_1 t \frac{d\eta}{d\xi}\right)_{\xi=-0.57} + \left(N_1 t \frac{d\eta}{d\xi}\right)_{\xi=0.57}\right) = 2\left(N_1 t \frac{d\eta}{d\xi}\right)_{\xi=-0.57}$$

Since $\xi = -1 + \frac{2\eta}{3a}$, we have that $\frac{d\eta}{d\xi} = \frac{3a}{2}$ and $\eta(\xi = -0.57) = \frac{3a}{2}0.43$.

$$2\left(N_1 t \frac{d\eta}{d\xi}\right)_{\xi=-0.57} = 3a[b\ c]^T (a - \frac{3a}{2}0.43)(2a - \frac{3a}{2}0.43)(3a - \frac{3a}{2}0.43)/(6a^3)$$

Problem 3

a) Premultiply and postmultiply $\boldsymbol{K} + \tilde{\boldsymbol{K}}$ by \boldsymbol{a}^T and \boldsymbol{a} , respectively. Identifying $\boldsymbol{N}\boldsymbol{a} = T$ and $\boldsymbol{B}\boldsymbol{a} = \boldsymbol{\nabla}T$ reveals that $\boldsymbol{a}^T(\boldsymbol{K} + \tilde{\boldsymbol{K}})\boldsymbol{a} > 0$ for $\alpha > 0$.

b) An example of an approximation is $T = \alpha_1 + \alpha_2 x + \alpha_3 y + \alpha_4 z + \alpha_5 xy + \alpha_6 xz + \alpha_7 yz + \alpha_8 xyz$. Show that the approximation satisfies the convergence criterion, i.e. completeness+compatibility requirements. See page 91-94.

Problem 4

The displacement, \boldsymbol{u} , is discretized as

$$\boldsymbol{u}(x,y) = \boldsymbol{N}(x,y)\boldsymbol{a} \tag{1}$$

where N and a are the shape functions and nodal values associated with the mechanical problem. The total strain is

$$\boldsymbol{\varepsilon} = \tilde{\boldsymbol{\nabla}} \boldsymbol{u} = \tilde{\boldsymbol{\nabla}} \boldsymbol{N}(x, y) \boldsymbol{a} = \boldsymbol{B} \boldsymbol{a} \quad \text{where} \quad \tilde{\boldsymbol{\nabla}} = \begin{bmatrix} \frac{\partial}{\partial x} & 0\\ 0 & \frac{\partial}{\partial y}\\ \frac{\partial}{\partial y} & \frac{\partial}{\partial x} \end{bmatrix}$$
(2)

The stress is calculated using Hooke's law, i.e.

$$\boldsymbol{\sigma} = \boldsymbol{D}\boldsymbol{\varepsilon}^{e} = \boldsymbol{D}(\boldsymbol{\varepsilon} - \boldsymbol{\varepsilon}^{0}) = \boldsymbol{D}(\boldsymbol{\varepsilon} - \alpha\Delta T(x, y) \begin{bmatrix} 1 & 1 & 0 \end{bmatrix}^{T})$$

The temperature difference is $\Delta T = T(x, y) - T_0$, where T_0 is the stress free temperature and the current temperature, T, is discretized as

$$T(x,y) = N^0 a^0$$

where N^0 and a^0 are the shape functions and nodal values associated with the temperature problem.

Problem 5

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	x	x	x			x	x	x			x	x	x		
K =	x	x	x			x	x	x			x	x	x		
	x	x	x	x	x	x	x	x	x	x	x	x	x	x	x
			x	x	x			x	x	x			x	x	x
			x	x	x			x	x	x			x	x	x
	x	x	x			x	x	x			x	x	x		
	x	x	x			x	x	x			x	x	x		
	x	x	x	x	x	x	x	x	x	x	x	x	x	x	x
			x	x	x			x	x	x			x	x	x
			x	x	x			x	x	x			x	x	x
	x	x	x			x	x	x			x	x	x		
	x	x	x			x	x	x			x	x	x		
	x	x	x	x	x	x	x	x	x	x	x	x	x	x	x
			x	x	x			x	x	x			x	x	x
			x	x	x			x	x	x			x	x	x





Problem 6 FHLF01

a)

One possible ansatz with the unknowns a_1 and a_2 which fulfils the boundary conditions is,

$$\psi = \left(\sin(\pi x)\sin(\pi y)\right)a_1 + \left(\sin(2\pi x)\sin(2\pi y)\right)a_2$$

This can also be written as

$$\psi = \left[\sin(\pi x)\sin(\pi y), \sin(2\pi x)\sin(2\pi y)\right] \left[\begin{array}{c} a_1\\ a_2 \end{array}\right] = \left[N_1, N_2\right] \left[\begin{array}{c} a_1\\ a_2 \end{array}\right] = \mathbf{N}\mathbf{a}$$

b)

Define

$$oldsymbol{B} =
abla oldsymbol{N} = \left[egin{array}{c} rac{dN_1}{dx}, & rac{dN_2}{dx} \ rac{dN_1}{dy}, & rac{dN_2}{dy} \end{array}
ight]$$

Stationarity of the functional then provides

$$\begin{split} & \frac{d}{d\epsilon} \Pi(\boldsymbol{a} + \epsilon \boldsymbol{c}) \big|_{\epsilon=0} = \\ & \frac{d}{d\epsilon} \left(\int_{\Omega} \frac{1}{2} (\boldsymbol{a} + \epsilon \boldsymbol{c})^T \boldsymbol{B}^T \boldsymbol{B} \right) (\boldsymbol{a} + \epsilon \boldsymbol{c}) - (\boldsymbol{a} + \epsilon \boldsymbol{c})^T \boldsymbol{N}^T [2(x+y) - 4] dV \right) \big|_{\epsilon=0} = \\ & \boldsymbol{c}^T \left(\left(\int_{\Omega} \boldsymbol{B}^T \boldsymbol{B} dV \right) \boldsymbol{a} - \int_{\Omega} \boldsymbol{N}^T [2(x+y) - 4] dV \right) \end{split}$$

Since this expression should be true for arbitrary variations $\boldsymbol{c},$ then the following equation holds,

$$Ka = f$$

where

$$\boldsymbol{K} = \int_{\Omega} \boldsymbol{B}^T \boldsymbol{B} dV \quad \boldsymbol{f} = \int_{\Omega} \boldsymbol{N}^T [2(x+y) - 4] dV$$

Problem 6 (FHL064) : (10p)

See Chapter 17 in Ottosen.